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Grüneisen parameter calculated from the single and and the polycrystal TOEC with that calculated bulk data. Using values for the linear expansion which,  $\alpha = 7.02 \times 10^{-6}$  °C<sup>-1</sup>, and specific heat, -0.005 cal g<sup>-1</sup> °C<sup>-1</sup> obtained from handbooks and density  $\rho = 8.578$  g/cm³, and isothermal bulk modified from the present in the relation

$$\gamma_B = 3\alpha K^T / C_v \rho \tag{4}$$

the value for the bulk Grüneisen parameter of -1.52. Using the methods of Ref. 4 for obtaining Grüneisen parameter from the elastic constants by eraging the contribution of 39 pure-mode phonons that parameter results in values of  $\gamma_{\rm SC}=1.511\pm0.026$  for the single crystal and the herystal, respectively. A useful calculational equation the polycrystal Grüneisen parameter which can obtained from the equations of Brugger<sup>4</sup> by imposing copy conditions is

$$+(K^T m_1' - \Delta K)/6C_{11}^S + (K^T m_2' - \Delta K)/3C_{44}^S,$$
 (5)

the slopes for the polycrystal and  $\Delta K = K^S - K^T$  is the elerence between the adiabatic and isothermal bulk stuli.

A second check on the values of the TOEC deterned here is to compare the measured polycrystal astants with values calculated from the singlestal constants. Recently, equations permitting this

comparison were derived<sup>30</sup> using strain-energy density considerations with the approximation that a uniform state of strain acting on the surface of a homogeneous, quasi-isotropic, polycrystalline body produces a uniform strain throughout the body. This development, analogous to that of Voigt relating the second-order elastic constants,<sup>31</sup> leads to the following relations between the TOEC:

$$\nu_1 = \frac{1}{3.5} (C_{111} + 18C_{112} + 16C_{123} - 30C_{144} - 12C_{166} + 16C_{456})$$

$$\nu_2 = \frac{1}{35} (C_{111} + 4C_{112} - 5C_{123} + 19C_{144} + 2C_{166} - 12C_{456})$$

$$\nu_3 = \frac{1}{35} (C_{111} - 3C_{112} + 2C_{123} - 9C_{144} + 9C_{166} + 9C_{456}). \tag{6}$$

Using the values of  $C_{ijk}$  for columbium given in Table III in these equations results in  $\nu_1 = -4.59 \pm 0.38$ ,  $\nu_2 = -3.80 \pm 0.15$ , and  $\nu_3 = +0.78 \pm 0.09$ , all in units of  $10^{12}$  dyn/cm<sup>2</sup>. These values are seen to be in very good agreement with the values measured for the polycrystal sample B shown in Table V which provides an indirect check on both sets of TOEC.

## ACKNOWLEDGMENTS

We are grateful for the excellent work in sample preparation done by Robert G. Herron who polished the samples and Paul C. Romo who performed the x-ray orientations.

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## Lithium Niobate: Effects of Composition on the Refractive Indices and Optical Second-Harmonic Generation\*

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Crystals of lithium niobate have been grown with varying amounts of Li<sub>2</sub>O, Nb<sub>2</sub>O<sub>3</sub>, and MgO present in the melt. The refractive indices have been measured and the birefringence so obtained correlated with the phase-matching temperature for second-harmonic generation (SHG) and with the composition. The use of optical methods is decribed for the study of compositional uniformity and a method of correcting the effects of uniform composition gradients on SHG described.

## I. INTRODUCTION

An extensive literature already exists on lithium bate with reference to its use in nonlinear optics for axing or modulating optical beams. It was first proced in this connection by Boyd *et al.*, who described

able for calculating the details of phase matching. Miller et al.<sup>2</sup> extended this data by pointing out the favorable property of thermally tunable refractive indices present in lithium niobate and showed how this could be used to obtain noncritical phase matching. Hobden and Warner<sup>3</sup> have given extensive details of the wavelength and thermal variation of the refractive

its nonlinear properties and refractive-index data suit-

Based in part on a dissertation submitted to the University London in partial fulfullment of the requirements for the degree Doctor of Philosophy.

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